

ECE3084-L10 Laplace Transform

Thursday, March 2, 2017 4:05 AM

Recall:

$$\text{Fourier Transform: } X(j\omega) = \int_{-\infty}^{\infty} x(t) e^{-j\omega t} dt$$

$$\underbrace{x_1(t) * x_2(t)}_{\text{convolution}} \iff \underbrace{X_1(j\omega) X_2(j\omega)}_{\text{multiplication}}$$

Particularly useful for freq domain analysis:

$$x(t) \longrightarrow \boxed{h(t)} \longrightarrow y(t)$$
$$Y(j\omega) = X(j\omega)H(j\omega)$$

$$x(t) = A \cos(\omega t + \phi)$$

$$\text{then } y(t) = |H(j\omega)| A \cos(\omega t + \phi + \angle H(j\omega))$$

The FT doesn't work for many signals, e.g.

$$x(t) = u(t)$$

$$X(j\omega) = \int_{-\infty}^{\infty} u(t) e^{-j\omega t} dt = \int_0^{\infty} e^{-j\omega t} dt \quad \leftarrow \text{This integral does not exist.}$$

However, the FT does exist if the step fn decays: $x(t) = e^{-at} u(t)$

If we can introduce a factor $e^{-\sigma t}$ into the transform, then it could work

for this input:

$$\begin{aligned} X(\sigma + j\omega) &= \int_0^{\infty} e^{-\sigma t} e^{-j\omega t} dt \quad \text{for } \sigma > 0 \\ &= \int_0^{\infty} e^{-(\sigma + j\omega)t} dt \\ &= \frac{1}{-(\sigma + j\omega)} e^{-(\sigma + j\omega)t} \bigg|_0^{\infty} \\ &= \frac{1}{\sigma + j\omega} \end{aligned}$$

Laplace Transform

Let $s = \sigma + j\omega$ $s \in \mathbb{C} \leftarrow$ set of complex numbers

$$X(s) = \int_{-\infty}^{\infty} x(t) e^{-st} dt \quad \text{Bilateral (or two-sided) Laplace Transform}$$

↑ We just replaced the $j\omega$ from the FT with a more general complex number, $s = \sigma + j\omega$

The FT projects signals onto complex sinusoids $e^{j\omega t}$, which are unable to represent some types of signals (e.g. signals that go to ∞ as $t \rightarrow \infty$).

The LT projects signals onto $e^{st} = e^{(\sigma + j\omega)t} = e^{\sigma t} e^{j\omega t}$.

These are complex sinusoids that can exponentially grow or decay in magnitude depending on what the value of σ is.

This is a more general basis set that allows us to capture a richer set of signals.

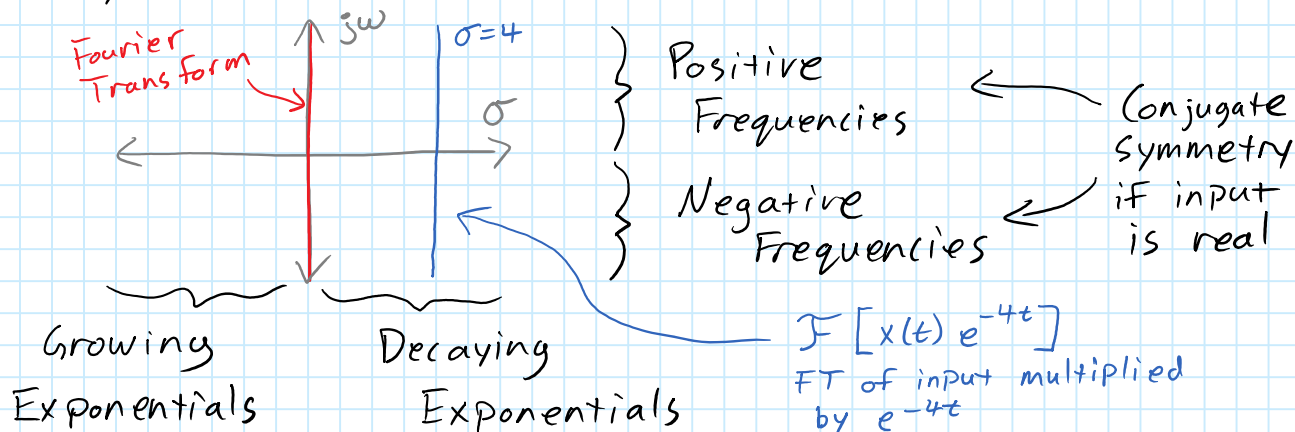
*Note that if we choose $\sigma = 0$ so that $s = j\omega$, then the LT reduces to the FT. In other words, the FT is a special case or subset of the LT.

For a particular signal $x(t)$, the region of convergence (ROC) is the set of all s for which the integral converges.

The ROC only constrains the real part of s (It's bounded by vertical lines in the s -plane).

One way to think about the LT:*

To find the values along a vertical line $s = \sigma$ in the s -plane, multiply the input by $e^{-\sigma t}$, then take the Fourier Transform.



*The LT is not really a two-stage process like this though.

The bilateral LT is rarely used by practicing engineers and complicates dealing with ROCs. For causal systems with inputs for $t \geq 0$ and some set of initial conditions at $t = 0^-$, we use:

Unilateral Laplace Transform ("Laplace Transform"):

$$\mathcal{L}[x(t)] = X(s) = \int_{0^-}^{\infty} x(t) e^{-st} dt$$

↑ Includes the point $t=0$ (ie if there's a δ there)

Unless we specifically say "bilateral," this is what we are referring to.

The system outputs $y(t)$ we derive with this are only valid for $t \geq 0$.

Examples:

$$\begin{aligned} 1. \quad x(t) = u(t) \quad X(s) &= \int_0^{\infty} u(t) e^{-st} dt \\ &= -\frac{1}{s} e^{-st} \Big|_0^{\infty} \\ &= -\frac{1}{s} e^{-s(\infty)} + \frac{1}{s} \end{aligned}$$

↑ goes to 0 if $\text{Re}(s) > 0$

$$X(s) = \frac{1}{s}$$

$$\boxed{\mathcal{L}[u(t)] = \frac{1}{s} \quad \text{Re}(s) > 0}$$

ROC is the RHP (right half plane)

$$\begin{aligned} 2. \quad x(t) = e^{-at} u(t) \quad X(s) &= \int_0^{\infty} e^{-at} e^{-st} dt \\ &= -\frac{1}{s+a} e^{-(s+a)t} \Big|_0^{\infty} \\ &= -\frac{1}{s+a} e^{-(s+a)\infty} + \frac{1}{s+a} \end{aligned}$$

↑ goes to zero if $\text{Re}(s) > -a$

$$\boxed{\mathcal{L}[e^{-at} u(t)] = \frac{1}{s+a} \quad \text{Re}(s) > -a}$$

I think most of the entries on the table with an e^{-at} have a ROC of $\text{Re}(s) > -a$.

$$3. \quad x(t) = t e^{-at}$$

$$X(s) = \int_0^{\infty} t e^{-at} e^{-st} dt$$

$$= \int_0^{\infty} t e^{-(a+s)t} dt$$

$$u = t \quad dv = e^{-(a+s)t}$$

$$du = 1 \quad v = \frac{-1}{a+s} e^{-(a+s)t}$$

Integration by parts:

$$\int u dv = uv - \int v du$$

(can derive this by integrating the product rule of differentiation.)

$$X(s) = \left(\frac{-1}{a+s} t e^{-(a+s)t} - \int \frac{-1}{a+s} e^{-(a+s)t} dt \right) \Big|_0^{\infty} \quad \text{Re}(s) > a$$

$$= \left(\frac{-1}{a+s} t e^{-(a+s)t} - \frac{1}{(a+s)^2} e^{-(a+s)t} \right) \Big|_0^{\infty}$$

$$= (0 - 0) - \left(0 + \frac{-1}{(a+s)^2} \right)$$

$$\mathcal{L} [t e^{-at} u(t)] = \frac{1}{(a+s)^2} \quad \text{Re}(s) > a$$

↑ The $u(t)$ shows up in every entry of the table (except where redundant) to emphasize that this is the unilateral LT.

In general, better to use tables & properties rather than integrating.

Key Properties of the LT (see table for others)

- Linearity: $\mathcal{L} [a x_1 + b x_2] = a \mathcal{L} [x_1] + b \mathcal{L} [x_2]$
 $= a X_1(s) + b X_2(s)$

Follows easily from linearity of integrals

- Time Delay: $\mathcal{L} [x(t-T)] = e^{-sT} X(s)$

$$\mathcal{L} [x(t-T)] = \int_0^{\infty} x(t-T) e^{-st} dt$$

Change of vars:

$$\tau = t - T \quad t = \tau + T$$

$$dt = d\tau$$

$$= \int_{\tau+T=0}^{\tau+T=\infty} x(\tau) e^{-s(\tau+T)} d\tau$$

$$= e^{-sT} \int_{-T}^{\infty} x(\tau) e^{-s\tau} d\tau$$

↑ Assume $x(t) = 0$ for $t \in [-T, 0)$

$$= e^{-sT} \int_0^{\infty} x(\tau) e^{-s\tau} d\tau$$

$$= e^{-sT} X(s)$$

• Differentiation: $\mathcal{L}[x'(t)] = sX(s) - x(0^-)$

$$\mathcal{L}[x'(t)] = \int_0^{\infty} x'(t) e^{-st} dt \quad \leftarrow \text{small } x!$$

Integration by parts: $\int u dv = uv - \int v du$

$$u = e^{-st} \quad dv = x'(t) dt$$

$$du = -s e^{-st} dt \quad v = x(t)$$

$$= \left(e^{-st} x(t) - \int x(t) (-s e^{-st}) dt \right) \Big|_0^{\infty}$$

$$= \underbrace{x(t) e^{-st} \Big|_0^{\infty}}_{0 - x(0)} + s \underbrace{\int_0^{\infty} x(t) e^{-st} dt}_{X(s)}$$

$$= sX(s) - x(0^-)$$

$\leftarrow 0^-$ because that's how we defined the lower bound of the LT. We just get lazy & don't always write it.

Can apply repeatedly:

$$\begin{aligned} \mathcal{L}[x''(t)] &= s \mathcal{L}[x'(t)] - x(0^-) \\ &= s^2 X(s) - s x'(0^-) - x(0^-) \end{aligned}$$

Notice: The derivatives disappeared and became polynomials in the s -domain. Thus, differential equations in the time domain become algebraic equations in the s -domain!

This is why the LT is so useful for solving diff eqs.

• Integration:

$$\text{Let } v(t) = \begin{cases} 0 & t < 0 \\ \int_0^t x(\tau) d\tau & t \geq 0 \end{cases}$$

$$\text{So } x(t) = v'(t)$$

$$\mathcal{L}[v'(t)] = sV(s) - v(0)$$

$$\mathcal{L}[x(t)] = sV(s) \quad \leftarrow = 0 \text{ by definition}$$

$$X(s) = sV(s)$$

$$V(s) = \frac{1}{s} X(s)$$

$$\mathcal{L}\left[\int_0^t x(\tau) d\tau\right] = \frac{1}{s} X(s)$$

• Final Value Theorem: $\lim_{t \rightarrow \infty} x(t) = \lim_{s \rightarrow 0} sX(s)$

If the limit exists, and $x(t)$ & $x'(t)$ have LTS.

Proof:
$$\begin{aligned} \lim_{s \rightarrow 0} sX(s) &= \lim_{s \rightarrow 0} (sX(s) - x(0) + x(0)) \\ &= \lim_{s \rightarrow 0} (\mathcal{L}[x'(t)]) + x(0) \\ &= \lim_{s \rightarrow 0} \int_0^{\infty} x'(t) e^{-st} dt + x(0) \\ &= \int_0^{\infty} x'(t) \left(\lim_{s \rightarrow 0} e^{-st} \right) dt + x(0) \\ &= \int_0^{\infty} x'(t) dt + x(0) \\ &= (x(\infty) - x(0)) + x(0) \\ &= x(\infty) \end{aligned}$$

• Initial Value Theorem: $\lim_{t \rightarrow 0^+} x(t) = \lim_{s \rightarrow \infty} sX(s)$

Proof: Similar, just change the limit to $s \rightarrow \infty$ instead of $s \rightarrow 0$.

Examples

• $x(t) = \cos(\omega t) u(t)$

$= \frac{1}{2} (e^{j\omega t} + e^{-j\omega t}) u(t)$

$$\begin{aligned} X(s) &= \frac{1}{2} \left(\frac{1}{s-j\omega} + \frac{1}{s+j\omega} \right) \\ &= \frac{1}{2} \left(\frac{s+j\omega + s-j\omega}{s^2 + \omega^2} \right) = \frac{1}{2} \left(\frac{2s}{s^2 + \omega^2} \right) \\ &= \frac{s}{s^2 + \omega^2} \end{aligned}$$

$x(t) e^{at} \xleftrightarrow{\mathcal{L}} X(s-a)$

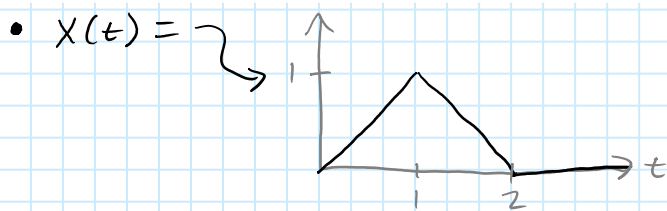
$u(t) \leftrightarrow \frac{1}{s}$

• $x(t) = t u(t) = \int_0^t u(\tau) d\tau$

$X(s) = \frac{1}{s} \mathcal{L}[u(t)] = \frac{1}{s^2}$

$\int_0^t x(\tau) d\tau \leftrightarrow \frac{1}{s} X(s)$

$u(t) \leftrightarrow \frac{1}{s}$



$$x(t) = t u(t) - 2(t-1)u(t-1) + (t-2)u(t-2)$$

$$X(s) = \frac{1}{s^2} - 2e^{-s} \left(\frac{1}{s^2} \right) + e^{-2s} \left(\frac{1}{s^2} \right)$$

$$= \frac{1}{s^2} (1 - 2e^{-s} + e^{-2s})$$

• $x(t) = \delta(t)$ $X(s) = \int_0^{\infty} \delta(t) e^{-st} dt = (e^{-st} |_{t=0}) = 1$

$$X(s) = 1$$

• $x(t) = \sin\left(t - \frac{\pi}{3}\right) u(t)$ Note: $\sin(A-B) = \sin A \cos B - \cos A \sin B$

$$= \left(\sin(t) \cos\left(\frac{\pi}{3}\right) - \cos(t) \sin\left(\frac{\pi}{3}\right) \right) u(t)$$

$$= \left(\frac{1}{2} \sin(t) - \frac{\sqrt{3}}{2} \cos(t) \right) u(t) \quad \omega=1$$

$$X(s) = \frac{1}{2} \left(\frac{1}{s^2+1} \right) - \frac{\sqrt{3}}{2} \left(\frac{s}{s^2+1} \right) = \frac{1}{2} \left(\frac{-\sqrt{3}s+1}{s^2+1} \right)$$

Why can't we use the time delay property?

→ the step fn was not delayed.